



## Investigations of GaN metal-oxide-semiconductor capacitors with sputtered HfO<sub>2</sub> gate dielectrics

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### ABSTRACT

This work elucidates the properties of Al/HfO<sub>2</sub>/GaN metal-oxide-semiconductor capacitors using reactively sputtered HfO<sub>2</sub> as a gate dielectric. The influence of GaN surface treatments and the post-annealing of HfO<sub>2</sub> films on the leakage current, flat-band voltage, interface trap densities, dielectric constants, and effective oxide charges of the GaN MOS capacitors are presented. The Ga oxynitride on the surface of GaN was effectively removed by chemical solutions that also slightly reduced the dielectric constant, slightly increased the flat-band voltage, eliminated the hysteresis of the capacitance–voltage measurement, and yielded a similar leakage to that without surface treatment. A highest dielectric constant of HfO<sub>2</sub> (17) was obtained when the sample was annealed at 600 °C for 20 min, while the lowest interface trap density ( $5.3 \times 10^{11} \text{ cm}^{-2}$ ) was obtained when the sample was annealed at 800 °C for 40 min. The leakage mechanism was well fitted by the Schottky emission and Frenkel–Poole emission models at a lower and higher electric field.

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III-Nitride semiconductors have been demonstrated to have potential applications in high-power, high-speed, and high-temperature devices because of their attractive fundamental properties such as high direct band-gap, chemical stability and high electron saturation velocity. Recently, GaN-based transistors such as metal-semiconductor field-effect-transistors (MES-FETs), hetero-junction FETs (HFETs), and high-electron-mobility transistors (HEMTs) have been reported [1–4]. All such devices suffer from the gate leakage problem. A large leakage current flows at the metal and GaN interface, making a good Schottky contact difficult to achieve. Metal-oxide-semiconductor FETs (MOS-FETs) and MOS-HFETs with an oxide interlayer between the metal and the semiconductor have much smaller leakage currents. Therefore, GaN-based MOS-FETs and MOS-HFETs have been important topics of research [5–8].

GaN-based MOS-FETs provide a number of advantages over Si-based devices. MOS capacitors are good forerunners of MOS-FETs or MOS-HFETs that can be employed to study the properties of the semiconductor–insulator interface. Previously, numerous gate dielectrics, such as Ga<sub>2</sub>O<sub>3</sub>, SiO<sub>2</sub>, Si<sub>3</sub>N<sub>4</sub>, Al<sub>2</sub>O<sub>3</sub>, HfO<sub>2</sub>, MgO, SiN<sub>x</sub>O<sub>y</sub>, Ta<sub>2</sub>O<sub>5</sub>, and ZrO<sub>2</sub> [9–16] have been studied using MOS structures. Among these, HfO<sub>2</sub> has attracted the most attention by far due to

its good thermal stability and high dielectric constant. Optimizing the GaN MOS structure involves two major challenges—lowering the interface traps and improving the oxide quality. Many effective techniques are available for depositing insulators on GaN. They include plasma-enhanced chemical vapor deposition (PE-CVD), e-beam evaporation, low-pressure CVD, photo-CVD, sputtering and atomic layer deposition [10–12,14,15]. Among these techniques, sputtering offers the advantages of simplicity process, low cost and high throughput. Besides, various chemical solutions such as H<sub>3</sub>PO<sub>4</sub>, H<sub>2</sub>SO<sub>4</sub> + H<sub>2</sub>O<sub>2</sub>, H<sub>3</sub>PO<sub>4</sub>, HF, and HCl have been applied to clean the GaN surface prior to the oxide deposition [9–13]. The performance of the MOS capacitors depends sensitively on the surface chemistry of as-grown GaN. However, the relationship between surface treatments of GaN and corresponding MOS properties has not been elucidated. Moreover, no research into GaN MOS capacitors with a sputtered HfO<sub>2</sub> as an oxide layer has been conducted.

This work systematically studies the Al/HfO<sub>2</sub>/GaN MOS capacitor using reactively sputtered HfO<sub>2</sub> as an insulator. The effect of GaN surface treatments and the annealing conditions of HfO<sub>2</sub> films on the leakage current, flat-band voltage, interface traps densities, dielectric constants and effective oxide charges of the GaN MOS capacitors are presented. The dominant leakage mechanism is also described.

### 1. Experimental

N-type GaN films were grown on (0001) sapphire substrate by metal-organic chemical vapor deposition (Aixtron rf-200/4). A low-temperature GaN nucleation

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layer, 25 nm thick, was first deposited at 520 °C. This was followed by a 1.5 μm GaN grown at 1050 °C. Although the threading dislocations of GaN could be reduced by increasing the thickness, the thickness of GaN was well chosen to yield a favorable surface flatness, which is key to a metal–semiconductor interface, with help of the *in situ* reflectometry using an 500 nm monochromatic light. The precursors were trimethylgallium (TMG) for Ga, ammonia for nitrogen and silane for n-type doping, with hydrogen carrier gas. The Si doping concentration was around  $2 \times 10^{18} \text{ cm}^{-3}$ . Before the oxide layers were deposited, the GaN samples were prepared by sequentially rinsing in de-ionized (DI) water, hot acetone, IPA, DI water, chemical solution (5 min), and finally repeating DI water/hot acetone/IPA/DI water rinsing. The chemical solutions adopted herein included diluted hydrochloric acid (HCl, H<sub>2</sub>O:HCl=1:1), buffer oxide etchant (H<sub>2</sub>O:HF=50:1) and potassium hydroxide (KOH, 1 M). HfO<sub>2</sub> thin films were deposited by reactive sputtering (dc, 200 W) using Hf metal as a target, in an Ar/O<sub>2</sub> ambient (Ar:O<sub>2</sub> = 10:5 sccm). The background and working pressure during sputtering were  $1 \times 10^{-6}$  and  $5 \times 10^{-3}$  Torr, respectively. The as-deposited HfO<sub>2</sub> was thermally annealed at temperatures from 500 to 800 °C, for 5, 20, 40, and 60 min. The circular 250 nm-thick Al gate electrode was formed by e-beam evaporation. The optimized post-annealing condition of the metal electrodes (10 min in nitrogen ambient at 400 °C) was used throughout. Binding energies was obtained by high-resolution X-ray photoemission spectroscopy (HRXPS, Ulvac/PHI Quantera SXM). The grazing-incident X-ray diffraction (GIXRD, Rigaku D/MAX2500) was used to characterize the phase of HfO<sub>2</sub> films. The electrical properties of the fabricated MOS capacitors were characterized by measuring high-frequency capacitance–voltage (C–V) and current density–voltage (J–V) curves with an Agilent 4294A precision impedance analyzer and an Agilent semiconductor parameter analyzer (E5270B mainframe with module E5287A), respectively.

## 2. Results and discussion

Quantitative analysis of the as-grown HfO<sub>2</sub> thin films by HR-XPS indicated that the compositional ratio of O/Hf approached 2 (not shown), even without thermal annealing, indicating a good stoichiometry manipulation of the sputtering. HR-XPS was also conducted to examine the HfO<sub>2</sub>/GaN interface. The GaN surfaces of the samples studied herein were treated with different aqueous solutions but all the HfO<sub>2</sub> thin films were deposited under identical conditions. *In situ* Ar<sup>+</sup> ion sputtering was used to remove slowly the HfO<sub>2</sub> until the underlying Ga 2p signals appeared during HR-XPS measurement. Fig. 1(a) and (b) plots corresponding data for O 1s and Ga 2p, respectively. The bottom profile is constructed from signals from the as-grown GaN for comparison. Both the O 1s and Ga 2p signals were shifted to a binding energy that was approximately 2 eV lower than that of untreated GaN. These shifts were independently of the chemical solutions selected. All of the broad Ga 2p<sub>1/2</sub> peaks in the vicinity of 1115 eV for the chemical treated samples could be deconvoluted into two peaks, exhibiting a chemical shift of ~2 eV, indicating a Ga oxynitride interfacial layer was presented. The origin of the interlayer is ascribed to the initial stages of sputtering because no such signal from the as-cleaned GaN surfaces could be resolved.

The effects of surface treatments on the electrical properties of GaN MOS diodes were studied by comparing the KOH, HCl, and BOE treated samples. An untreated sample was used for reference. The leakage currents for all of the treated and non-treated MOS samples were as small as  $\sim 2 \times 10^{-9} \text{ A/cm}^2$  (Fig. 2(a)). Surface treatments

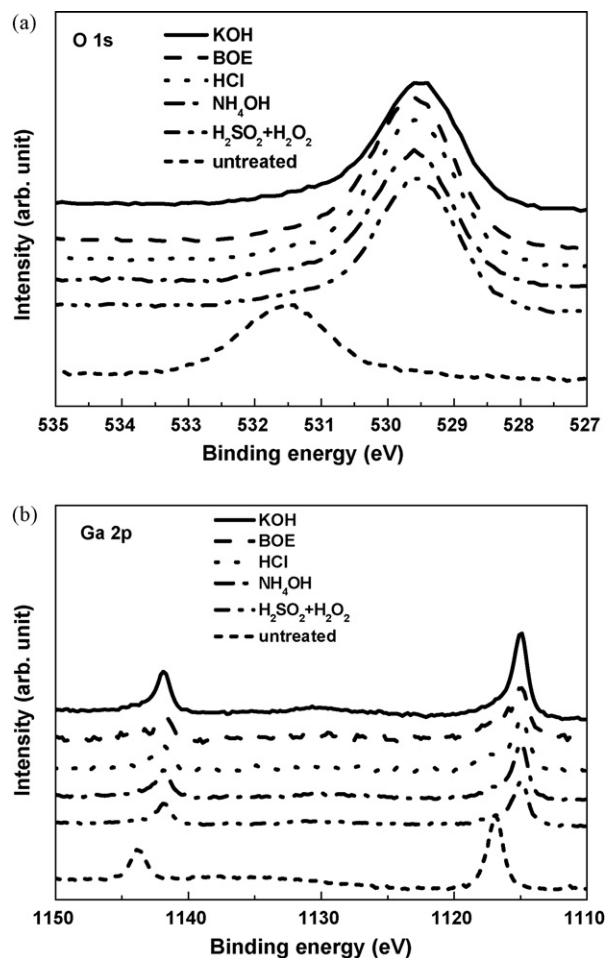


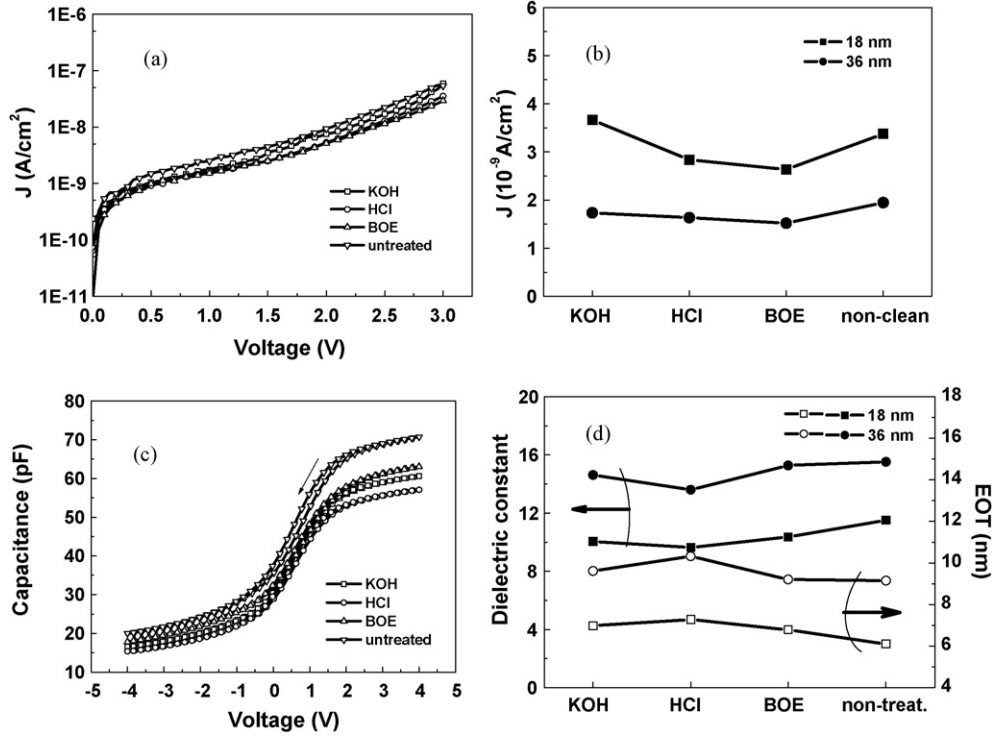
Fig. 1. HR-XPS profiles of the GaN/HfO<sub>2</sub> interfaces: (a) O 1s and (b) Ga 2p.

caused small deviations of the leakages. Regardless of treatment methods, the leakage current of MOS capacitor declined as the HfO<sub>2</sub> thickness increased (Fig. 2(b)), indicating that the main leakage current comes from the oxide layers not the interfaces. Fig. 2(c) plots well-behaved C–V curves that were measured at 1 MHz. No frequency dispersion was observed (100 kHz to 1 MHz), suggesting that the films have promising dielectric properties. Electrical hysteresis of the C–V curves due to the trapped charges within the oxide layers of the non-treated samples was observed, but decreased to a negligible level upon treatment with aqueous solution. Since the HfO<sub>2</sub> films were deposited under the same conditions, the hysteresis behavior was attributed to the native Ga oxynitride over the as-grown GaN surface. Moreover, the surface-treated samples had a lower accumulation capacitance than the non-treated one. From

Table 1

$\epsilon_{\text{GaN}}$ ,  $EOT$ ,  $V_{\text{FB}}$ ,  $N_{\text{eff}}$  and  $D_{\text{it}}$  values for Al/HfO<sub>2</sub>/GaN MOS capacitors with different annealing temperatures and times.

Temperature (°C)	Time (min)	$\epsilon_{\text{GaN}}$	$EOT$ (nm)	$V_{\text{FB}}$ (V)	$N_{\text{eff}}$ ( $\times 10^{12} \text{ cm}^{-2}$ )	$D_{\text{it}}$ ( $\times 10^{11} \text{ cm}^{-2}$ )
600	20	17	8.1	0.77	2.1	8.1
	40	14	10.1	0.75	1.6	6.8
	60	14	9.9	0.67	1.5	6.6
700	20	13		0.85	1.7	10.4
	40	11	12.6	0.86	1.4	9.5
	60	10	13.8	0.79	1.3	7.4
800	20	14	10.2	0.73	2.0	6.4
	40	12	11.2	0.72	1.5	5.3



**Fig. 2.** Electrical properties of the Al/HfO<sub>2</sub>/GaN MOS capacitors with different surface treatments of GaN and HfO<sub>2</sub> thickness. (a) *J*-*V* curves, (b) leakage currents at a +1 V gate bias, (c) high-frequency *C*-*V* curves and (d) dielectric constants and *EOT* values of HfO<sub>2</sub>. The *C*-*V* curves were measured from two sweep directions. The electrical hysteresis was presented for the non-treated sample.

the accumulation capacitance ( $C_{ox}$ ), the dielectric constants ( $\epsilon_{ox}$ ) are derived using:

$$C_{ox} = \frac{\epsilon_{ox}\epsilon_0 A}{t_{ox}}, \quad (1)$$

where  $\epsilon_0$  is the permittivity of a vacuum;  $t_{ox} = 18$  nm (or 36 nm) is the oxide thickness, determined by transmission electron microscopy (not shown), and  $A = 2 \times 10^{-4}$  cm<sup>2</sup> is the area of the capacitors. The equivalent oxide thickness (*EOT*), without deducting the quantum mechanical effect is defined as

$$EOT = \frac{\epsilon_{SiO_2}}{\epsilon_{ox}} t_{ox}, \quad (2)$$

where  $\epsilon_{SiO_2} = 3.9$  is the dielectric constant of SiO<sub>2</sub>. Fig. 2(d) plots the calculated dielectric constants of HfO<sub>2</sub> and *EOT* for the various surface treatment methods. The similarity among the degrees of leakage reveals that these samples had similar dielectric constants and *EOT*. Notably, thicker oxides are associated with larger dielectric constants (Fig. 2(d)), perhaps because of the improvement of the crystallinity. The capacitance of the n-GaN surface under the flat-band condition is expressed as

$$C_{FBS} = \frac{\epsilon_{GaN}\epsilon_0}{\lambda_n}, \quad (3)$$

where  $\lambda_n = (\epsilon_{GaN}\epsilon_0 kT/e^2 N_D)^{1/2}$  is the extrinsic Debye length;  $\epsilon_{GaN} = 8.9$  is the dielectric constant of GaN;  $k$  is the Boltzmann constant;  $T$  is the absolute temperature;  $e$  is the electronic charge, and  $N_D = 2 \times 10^{18}$  cm<sup>-3</sup> is the doping concentration. If the surface states of GaN are ignored, the flat-band capacitance,  $C_{FBS}$ , can be calculated by substituting  $C_{FBS}$  from Eq. (3) into the series capacitance relationship:

$$\frac{1}{C_{FB}} = \frac{1}{C_{ox}} + \frac{1}{C_{FBS}}. \quad (4)$$

Thus, the flat-band voltage ( $V_{FB}$ ) is obtained from the high-frequency *C*-*V* curves (Fig. 2(c)), plotted in Fig. 3(a). The difference between the metal-semiconductor work functions ( $\phi_{Ms}$ ) for Al metal and n-GaN is about  $-0.05$  eV by calculation, indicating that the MOS diodes can be under flat-band condition at near zero gate bias. The deviation of  $\phi_{Ms}$  and  $V_{FB}$  yields the effective oxide charges,  $Q_{eff} = Q_{ot} + Q_m + Q_f$ , using

$$Q_{eff} = \frac{C_{ox}(\phi_{ms} - V_{FB})}{A}, \quad (5)$$

where  $Q_{ot}$  denotes the oxide trapped charge density;  $Q_m$  is the mobile ionic charge density, and  $Q_f$  is the fixed charge density. Dividing the  $Q_{eff}$  by  $e = 1.6 \times 10^{-19}$  C yields the effective number density ( $N_{eff}$ ) of the oxides. From the normalized  $C_{FB}/C$  versus  $V$  curves yield the interface state density ( $D_{it}$ ) was calculated according to the well-known Terman method for MOS structures [17]. Fig. 3(a)–(c) plots the results of comparative studies of  $V_{FB}$ ,  $N_{eff}$  and  $D_{it}$  of the MOS samples, with and without GaN surface treatments. Fig. 3(a) reveals that the  $V_{FB}$  shifted positively as HfO<sub>2</sub> thickness increased. The dependence of  $V_{FB}$  on thickness revealed that the fixed charges were uniformly located throughout the bulk HfO<sub>2</sub>. Since the shift of  $V_{FB}$  came mostly from the oxide charges, the  $V_{FB}$  values of various treated samples were similar. The calculated  $N_{eff}$  in Fig. 3(b) confirmed this phenomenon. Fig. 3(c) indicates that the density of the interface charges of non-treated sample was lower than that of the chemically treated samples. We suggest two explanations. (1) The oxides on the GaN surface were removed using chemical solutions that also damaged the GaN surface. Thus, the surface defects contributed to the  $D_{it}$  of the GaN/HfO<sub>2</sub> interfaces. However, this possibility was ruled out because HR-XPS revealed no obvious chemical shift of the as-treated GaN surface. (2) A thin interfacial layer formed in the initial deposition of HfO<sub>2</sub>. According to Fig. 2(c) and based on the series capacitance, this interfacial layer may have a lower dielectric constant but fewer oxide traps than the

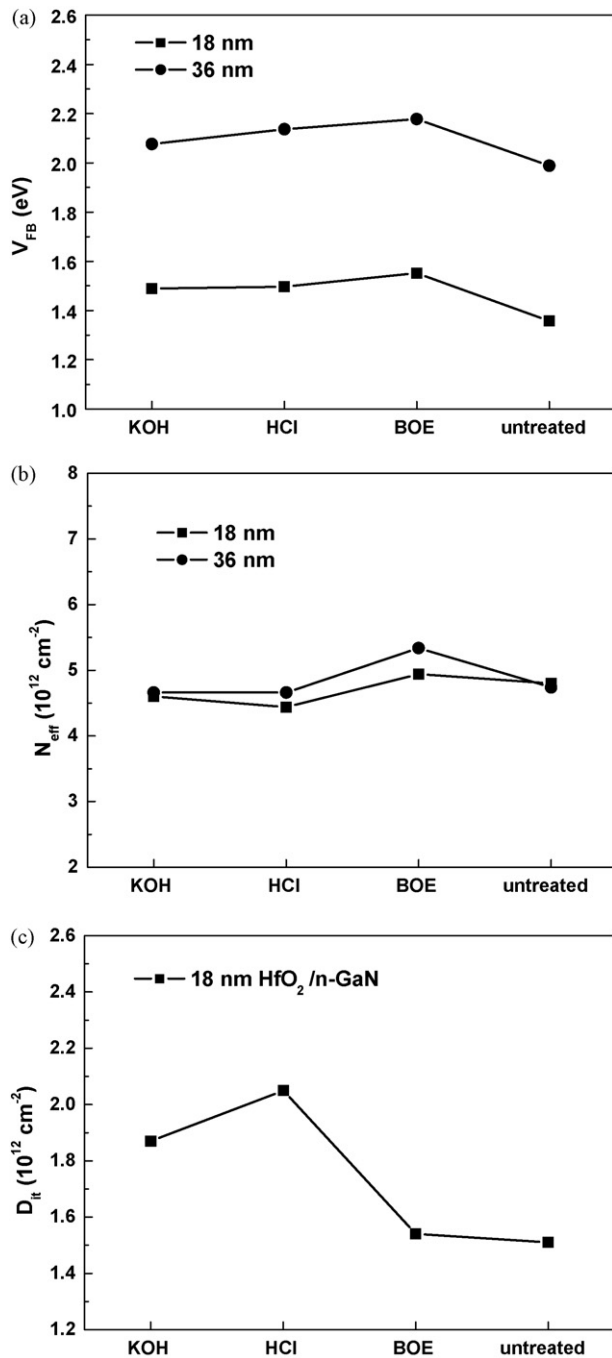


Fig. 3. (a) Flat-band voltage ( $V_{FB}$ ), (b) effective numbers density of oxide ( $N_{eff}$ ) and (c) interface trap densities for the GaN MOS capacitors with various surface treatments. The thickness of HfO<sub>2</sub> films were 18 and 36 nm.

native Ga oxynitride layer. Moreover, the Ga oxynitride layer of the as-grown GaN may act as a barrier to the radiation damage. The latter reason suggests that a buffer layer between the sputtered HfO<sub>2</sub> and the chemically treated n-GaN may be essential for further improvement.

Fig. 4 displays the GIXRD patterns of the 18 nm-thick-HfO<sub>2</sub> thin films annealed at elevated temperatures for 30 min. Monoclinic-related peaks became stronger and sharper as the temperature increased. An obvious transition from amorphous-like to monoclinic occurred when the annealing temperature exceeded 600 °C. Fig. 5 plots the leakage current density in the MOS capacitors that

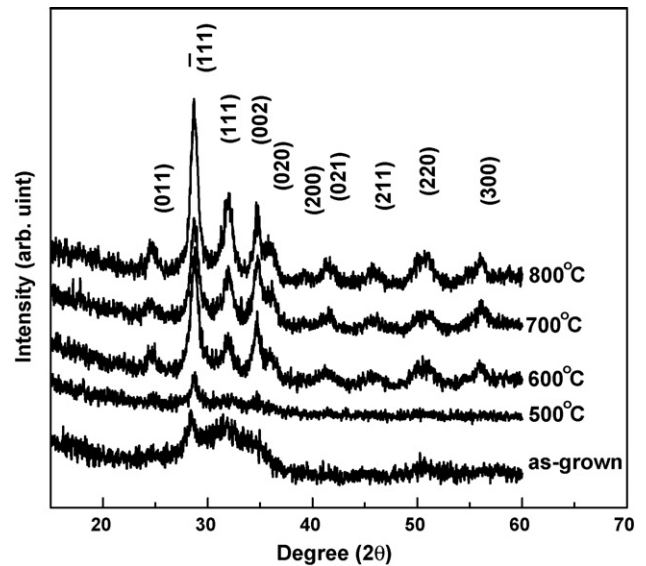


Fig. 4. GIXRD patterns for the HfO<sub>2</sub> annealed at temperatures from 500 to 800 °C.

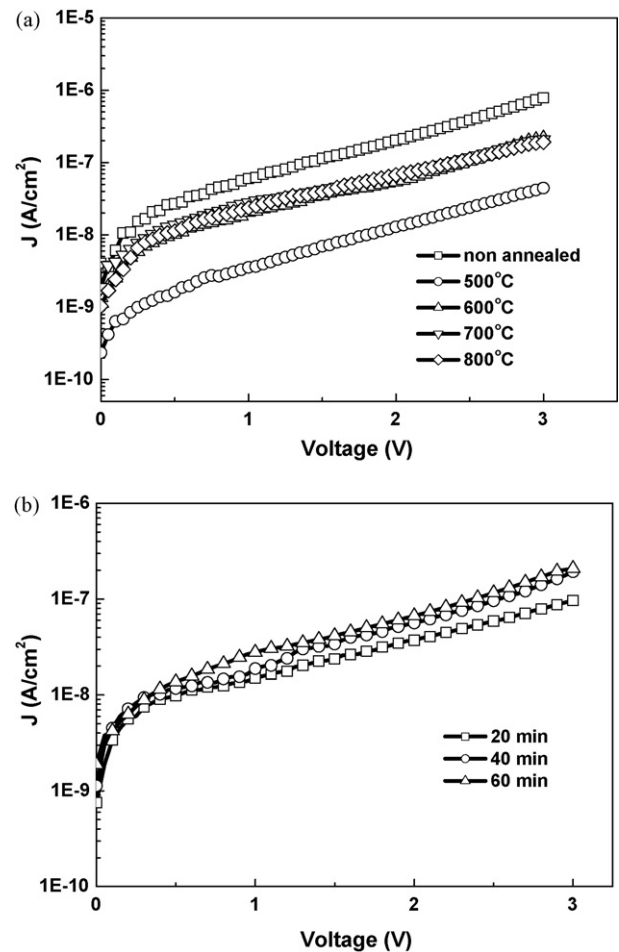
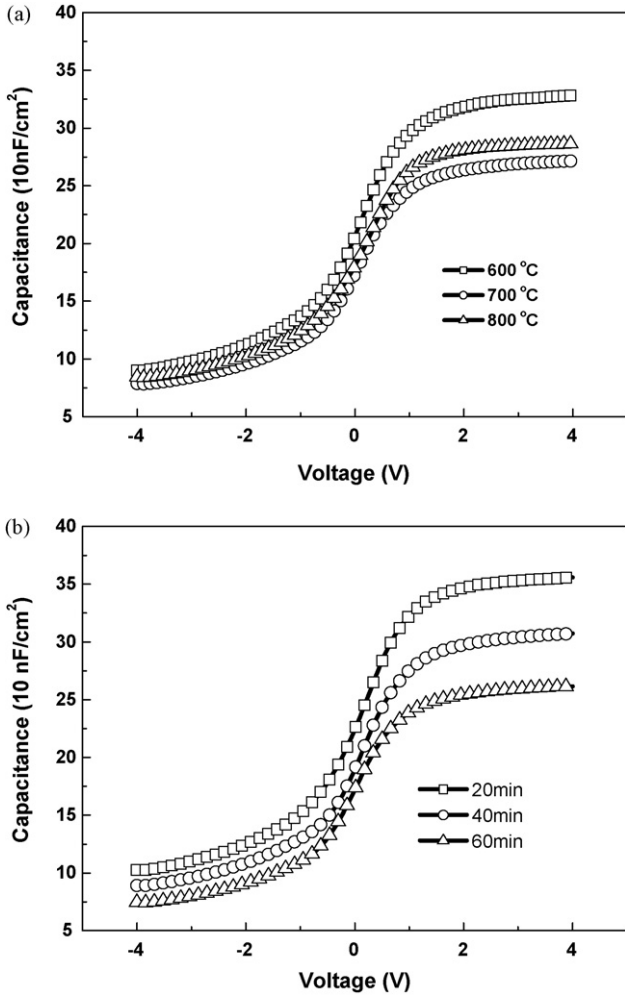


Fig. 5.  $J$ - $V$  curves of GaN MOS capacitors with HfO<sub>2</sub> annealed (a) at 500–800 °C for 40 min and (b) at 700 °C for 20–60 min. All of the samples were annealed in N<sub>2</sub> ambient.





**Fig. 6.** C–V curves of the GaN MOS capacitors with HfO<sub>2</sub> annealed (a) at 500–800 °C for 40 min and (b) at 700 °C for 20–60 min. All of the samples were annealed in N<sub>2</sub> ambient.

were annealed at various temperatures as a function of time. The leakage current density at a +1 V gate bias was markedly reduced from  $7 \times 10^{-8}$  to  $3 \times 10^{-9}$  A/cm<sup>2</sup> upon annealing at 500 °C (Fig. 5(a)). The leakage current density increased to about  $2 \times 10^{-8}$  A/cm<sup>2</sup> as the annealing temperatures increased (600–800 °C), because of grain boundaries were formed by high-temperature annealing. Changes in the annealing time at 700 °C in N<sub>2</sub> ambient had little effect on the leakage behaviors (Fig. 5(b)). Fig. 6(a) and (b) presents the C–V profiles of the MOS with different post-annealing temperatures and times. The oxide capacitance decreased as the annealing time increased. The highest dielectric constant, lowest  $N_{eff}$ , and lowest  $D_{it}$  within the experiment were found to be 17 (600 °C, 20 min),  $1.3 \times 10^{12}$  cm<sup>-2</sup> (700 °C, 60 min) and  $5.3 \times 10^{11}$  (800 °C, 40 min), respectively. Table 1 presents more  $\epsilon_{GaN}$ ,  $EOT$ ,  $V_{FB}$ ,  $N_{eff}$  and  $D_{it}$  values of the Al/HfO<sub>2</sub>/GaN MOS capacitors with different annealing temperatures and times. These values decreased as the annealing time increased.

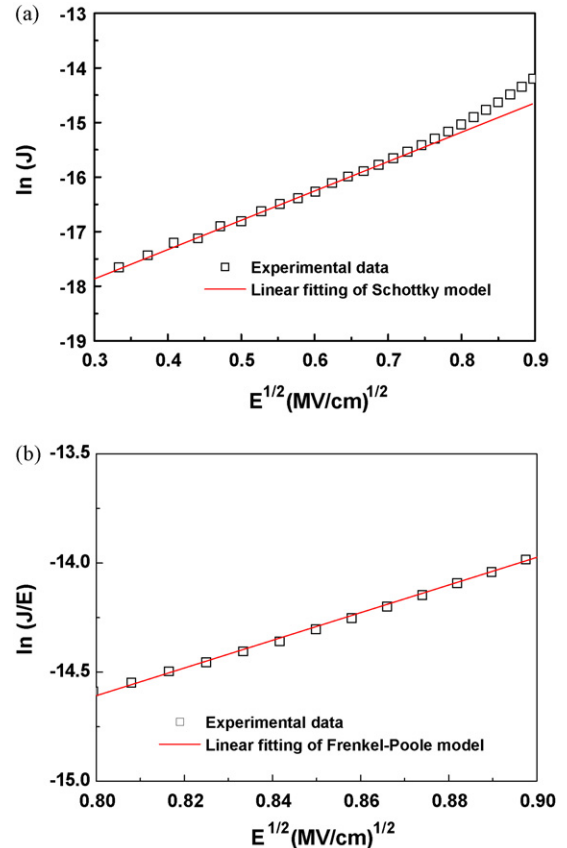
The  $J$ – $V$  characterization of MOS capacitor with moderate electrical properties (HfO<sub>2</sub> annealed at 700 °C, 40 min) was studied to understand further the leakage behaviors. The Schottky emission mechanism [18] is described by

$$J = A^* T^2 \exp \left[ \frac{-e(\phi_B - \sqrt{eE/4\pi\epsilon_r\epsilon_0})}{kT} \right], \quad (6)$$

where the effective Richardson constant  $A^*$  is a function of electron effective mass in oxide;  $J$  is the current density;  $k$  is the Boltzmann constant;  $h$  is the Planck constant;  $E$  is the electric field;  $\epsilon_r$  is the dynamic dielectric constant, and  $\phi_B$  is the Schottky barrier height. This equation was used to fit the leakage current behaviors in the low electric field side (0.3–0.75 MV/cm) of the  $\ln J$  versus  $\sqrt{E}$  plot. The leakage current in an electric field in the range of 0.3–0.75 MV/cm was fitted well with the Schottky emission mechanism (Fig. 7(a)). The  $\ln J$  versus  $\sqrt{E}$  plot also yields  $\phi_B$  and  $\epsilon_r$  as 1.12 eV and 9.3, respectively. The  $\epsilon_r$  calculated from this model was close to, but a little lower than the high-frequency C–V measurement. Notably, the value of  $\phi_B$  was very close to the conduction band offset ( $\sim 1.1$  eV) at the GaN/HfO<sub>2</sub> interface that was calculated using methods based on the charge neutrality [19]. When the electric field exceeded 0.75 MV/cm, the  $\ln J$ – $\sqrt{E}$  plot positively diverged from the Schottky emission mode, indicating that additional leakage paths contributed to the leakage current. The  $\ln(J/E)$  versus  $\sqrt{E}$  curve, plotted in Fig. 7(b), are fitted well by the Frenkel–Poole (FP) emission model, which is given by [20]

$$J \propto E \exp \left[ \frac{-q(\phi_B - \sqrt{eE/4\pi\epsilon_r\epsilon_0})}{kT} \right]. \quad (7)$$

The leakage characteristics indicate that the electron emission from the defect states within the oxide layer started when the electric field exceeded their activation energies.



**Fig. 7.** (a) Schottky and (b) Frenkel–Poole emission models fitted for the leakage behaviors of GaN MOS capacitors.

### 3. Conclusion

HfO<sub>2</sub> films were deposited on the n-type GaN by a reactive sputtering. The Al/HfO<sub>2</sub>/GaN MOS capacitors using the sputtered HfO<sub>2</sub> gate dielectric were investigated. The influence of GaN surface treatments and the post-annealing of HfO<sub>2</sub> films on the leakage current, flat-band voltage, interface trap densities, dielectric constants, and effective oxide charges of the GaN MOS capacitors were presented. The Ga oxynitride on the surface of GaN was effectively removed by chemical solutions that also slightly reduced the dielectric constant, slightly increased the flat-band voltage, eliminated the hysteresis of the capacitance–voltage measurement, and yielded a similar leakage to that without surface treatment. BOE treated sample showed lowest interface trap density to the other methods. A highest dielectric constant of HfO<sub>2</sub> (17) was obtained when the sample was annealed at 600 °C for 20 min, while the lowest interface trap density ( $5.3 \times 10^{11} \text{ cm}^{-2}$ ) was obtained when the sample was annealed at 800 °C for 40 min. The  $\epsilon_{\text{GaN}}$ ,  $E_{\text{OT}}$ ,  $V_{\text{FB}}$ ,  $N_{\text{eff}}$ , and  $D_{\text{it}}$  decreased monotonically as the annealing time increased. The leakage mechanism was well fitted by the Schottky emission and Frenkel–Poole emission model at a lower and higher electric field.

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### References

- [1] A. Koudymov, G. Simin, M.A. Khan, A. Taraki, R. Gaska, M.S. Shur, IEEE Electron Device Lett. 24 (2003) 680.
- [2] C.Q. Chen, J.P. Zhang, V. Adivarahan, A. Koudymov, J. Fatima, G. Simin, J. Yang, M.A. Khan, Appl. Phys. Lett. 82 (2003) 4593.
- [3] Y.K. Su, S.C. Wei, R.L. Wang, S.J. Chang, C.H. Ko, T.M. Kuan, IEEE Electron Device Lett. 24 (2003) 622.
- [4] Y. Irokawa, B. Luo, F. Ren, C.C. Pan, G.T. Chen, J.I. Chyi, S.S. Park, Y.J. Park, S.J. Pearton, Electrochem. Solid-State Lett. 7 (2004) G8.
- [5] M. Asif Khan, X. Hu, A. Tarakji, G. Simin, J. Yang, R. Gaska, Ms.S. Shur, Appl. Phys. Lett. 77 (2000) 1339.
- [6] W. Huang, T. Khan, T.P. Chow, IEEE Electron Device Lett. 27 (2006) 796.
- [7] K. Matocha, T. Paul Chow, R.J. Gutmann, IEEE Trans. Electron Device 52 (2005).
- [8] P.D. Ye, B. Yang, K.K. Ng, J. Bude, G.D. Wilk, S. Halder, J.C.M. Hwang, Appl. Phys. Lett. 86 (2005) 063501.
- [9] C.T. Lee, H.W. Chen, H.Y. Lee, Appl. Phys. Lett. (2003) 4304.
- [10] K. Matocha, R.J. Gitmann, T. Paul Chow, IEEE Trans. Electron Device 50 (2005) 1200.
- [11] S. Arulkumaran, T. Egawa, H. Ishikawa, T. Jimbo, M. Umeno, Appl. Phys. Lett. 73 (2003) 809.
- [12] Y.C. Chang, H.C. Chiu, Y.J. Lee, M.L. Huang, K.Y. Lee, M. Hong, Y.N. Chiu, J. Kwo, Y.H. Wang, Appl. Phys. Lett. 90 (2007) 232904.
- [13] J. Kim, B. Gila, R. Mahandru, J.W. Johnson, J.H. Shin, K.P. Lee, B. Luo, A. Onstine, C.R. Abernathy, S.J. Pearton, Ren, J. Electrochem. Soc. 149 (2002) G482.
- [14] S.J. Chang, C.K. Wang, Y.K. Su, C.S. Chang, T.K. Lin, T.K. Ko, H.L. Liu, J. Electrochem. Soc. 152 (2005) G423.
- [15] L.W. Tu, W.C. Kuo, K.H. Lee, P.H. Tsao, C.M. Lai, A.K. Chu, J.K. Sheu, Appl. Phys. Lett. 77 (2000) 3788.
- [16] S. Abermann, G. Pozzovivo, J. Kuzmik, G. Strasser, D. Pogany, J.-F. Carlin, N. Grandjean, E. Bertagnolli, Semicond. Sci. Technol. 22 (2007) 1272.
- [17] L.M. Terman, Solid-State Electron. 5 (1962) 258.
- [18] T. Yamaguchi, H. Satake, N. Fukushima, A. Toriumi, Tech. Dig.: Int. Electron Devices Meet (2000) 19.
- [19] J. Robertson, B. Falabretti, J. Appl. Phys. 100 (2005) 014111.
- [20] J. Frenkel, Phys. Rev. 54 (1938) 647.